

A convergent discretization of the porous medium equation with fractional pressure

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Based on a joint work with:

- E. R. Jakobsen - NTNU (Trondheim, Norway)

- 1** Introduction
- 2 Numerical scheme (Part 1)
- 3 Numerical scheme (Part 2)

Porous Medium Equation / Fast Diffusion Equation

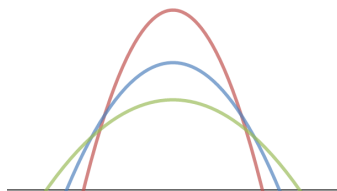
$$u_t - \Delta u^m = 0 \quad \text{in } \mathbb{R}^N \times (0, T) \quad (\text{PME/FDE})$$

Self Similar solutions: $\mathcal{U}(x, t) = t^{-\alpha} F(|x|t^{-\beta})$

Finite/Infinite speed of propagation with $m = 1$ as borderline cases

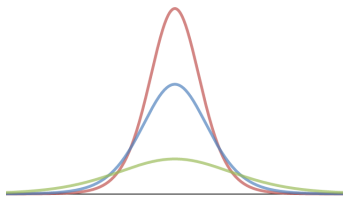
Slow Diffusion: $m > 1$

$$F(y) \sim (R^2 - |y|^2)_+^{1/(m-1)}$$



Fast Diffusion : $m < 1$

$$F(y) \sim (R^2 + |y|^2)^{-1/(1-m)}$$



Reference: *Vázquez's book 2007 on the PME*

Fractional porous medium models

$$u_t - \Delta u^m = 0$$

↓

$$u_t + (-\Delta)^s u^m = 0$$

de Pablo, Quirós, Rodríguez
and Vázquez

Infinite speed of
propagation for all $m > 0$.

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Caffarelli and Vázquez ($m=2$)
& Stan, del Teso, and Vázquez

$$m \in (1, \infty)$$

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Biler, Imbert,
Karch and Monneau

Finite speed of
propagation for all $m > 1$.

$$u_t = \nabla \cdot (u^{m-1} \nabla (-\Delta)^{-\sigma} u)$$

Continuity equation: $u_t = -\nabla \cdot (u \mathbf{v})$

$u \equiv$ density (mass per unit volume), $\mathbf{v} \equiv$ flow velocity field.

In this case

$$\mathbf{v} = u^{m-2} \nabla (-\Delta)^{-\sigma} u$$

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- The velocity field is nonlocal.

$$\nabla (-\Delta)^{-\sigma} f(x) = \nabla^{1-2\sigma} f(x) = \int (f(x) - f(y)) \frac{x-y}{|x-y|^{N+2-2\sigma}} dy$$

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Consequences:

- Finite speed of propagation if $m > 2$ (also $m = 2$, Caffarelli and Vázquez 11')
- Infinite speed of propagation if $m \in (1, 2)$.

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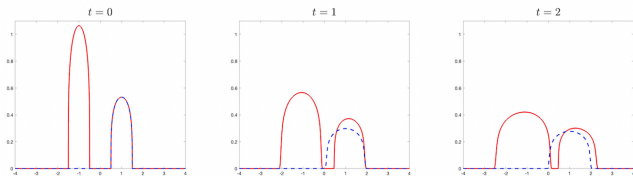
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Consequences:

- Finite speed of propagation if $m > 2$ (also $m = 2$, Caffarelli and Vázquez 11')
- Infinite speed of propagation if $m \in (1, 2)$.
- No comparison principle (Caffarelli and Vázquez 11')



$$\begin{aligned}\partial_t u - \nabla \cdot (u^{m-1} \nabla (-\Delta)^{-\sigma} u) &= 0, \\ u(\cdot, 0) &= \mu_0.\end{aligned}$$

Weak solution

$$\int_0^T \int_{\mathbb{R}^d} u \phi_t - \int_0^T \int_{\mathbb{R}^d} u^{m-1} \nabla (-\Delta)^{-\sigma} u \nabla \phi + \int_{\mathbb{R}^d} \phi(x, 0) d\mu_0(x) = 0,$$

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Properties (Stan, dT & Vázquez – 19')

- Existence of nonnegative weak solutions for $\mu_0 \in \mathcal{M}_+(\mathbb{R}^d)$.
- Conservation of mass.
- L^∞ -decay

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- (*dT, Jakobsen*) Tightness: For every $\varepsilon, T > 0$ there exists $R = R(\varepsilon, \mu_0, T) > 0$ s.t.

$$\sup_{t \in [0, T]} \int_{|x| > R} u(x, t) dx < \varepsilon.$$

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Integrating the equation

$$\partial_t v + (\partial_x v)^{m-1} (-\Delta)^s v = 0,$$

with $s = 1 - \sigma$.

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 - It is not in divergence form, and well-posedness has to be understood in the viscosity solution sense.
 - A general class of nonlocal quasilinear equations including this one was studied in (Chasseigne and Jakobsen 17'):

$$\begin{aligned} L[\phi](x) &= \int_{|z|>0} (\phi(x + f(\nabla\phi(x))z) - \phi(x)) \frac{dz}{|z|^{d+2s}} \\ &= -f(\nabla\phi(x))^{2s} (-\Delta)^s \phi(x) \end{aligned}$$

$$L[\phi] = (\partial_x \phi)^{m-1} (-\Delta)^s \phi$$

- Let us recall the definition of the p -Laplacian

$$\Delta_p \phi = \nabla \cdot (|\nabla \phi|^{p-2} \nabla \phi)$$

that in dimension $d = 1$ can be expressed as

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- (Chasseigne and Jakobsen 17') The operator L is a nonlocal version of the p -Laplacian. It does not coincide with other ones in the literature

$$\text{(Variational)} \quad (-\Delta)_p^s \phi(x) = \int \frac{|\phi(x) - \phi(x+z)|^{p-2} (\phi(x) - \phi(x+z))}{|z|^{d+sp}} dz$$

$$\text{(Game th.)} \quad -\Delta_p^s \phi(x) = \sup_{|y|=1} \int_{\mathcal{E}(x,y)} \frac{\phi(x) - \phi(x+z)}{|z|^{d+sp}} dz + \inf_{|\bar{y}|=1} \int_{\mathcal{E}(x,\bar{y})} \frac{\phi(x) - \phi(x+z)}{|z|^{d+sp}} dz$$

$$u_t - \partial_x (u^{m-1} \partial_x (-\partial_{xx})^{-\sigma} u) = 0 \quad (\text{FPE})$$

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- $u \geq 0$ (nonnegative weak sol.) $\iff \partial_x v \geq 0$ (nondecreasing viscosity sol.)

By (Chasseigne and Jakobsen 17'):

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By (Chasseigne and Jakobsen 17'):

- (IP) has existence and uniqueness of viscosity solutions.

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By (Chasseigne and Jakobsen 17'):

- (IP) has existence and uniqueness of viscosity solutions.
- (IP) has comparison principle for sub and super solutions.

This seems to be a right context for numerical schemes!

- 1 Introduction
- 2 Numerical scheme (Part 1)**
- 3 Numerical scheme (Part 2)

How to discretize

$$L[\phi] = (\partial_x \phi)^{m-1} (-\Delta)^s \phi$$

so that it preserves the monotonicity property:

If $\psi(x) = \phi(x)$ and $\psi \geq \phi$ then $L[\psi](x) \geq L[\phi](x)$.

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• *Discretization of the fractional Laplacian:*

$$-(-\Delta)^s \phi(x) = \text{P.V.} \int_{|z|>0} (\phi(x+z) - \phi(x)) \frac{dz}{|z|^{1+2s}}$$

$$-(-\Delta)_h^s \phi(x) \sim \sum_{k \neq 0} (\phi(x+z_k) - \phi(x)) \omega_k \quad \text{with} \quad \omega_k = \omega_{-k} \geq 0.$$

such that for $\phi \in C^2 \cap C_b$

$$(-\Delta)_h^s \phi \xrightarrow{h \rightarrow 0} (-\Delta)^s \phi \quad \text{in} \quad L_{\text{loc}}^\infty.$$

- *Discretization of ∂_x* : We think of $L[\phi] = (\partial_x \phi)^{m-1} (-\Delta)^s \phi$ as

$$L[\phi] = (\partial_x \phi)^{m-1} b(x, t),$$

with b changing signs.

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with b changing signs. Thus, we choose an UPWIND-type discretization:

$$\partial_x \phi(x) \sim D_h \phi(x) := \begin{cases} \frac{\phi(x+h) - \phi(x)}{h} & \text{if } (-\Delta)_h^s \phi(x) \leq 0, \\ \frac{\phi(x) - \phi(x-h)}{h} & \text{if } (-\Delta)_h^s \phi(x) > 0. \end{cases}$$

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- *Complete discretization*

$$L_h[\phi](x) = (D_h \phi(x))^{m-1} (-\Delta)_h^s \phi(x)$$

- (Consistent) For $\phi \in C^2 \cap C_b$, $L_h[\phi] \xrightarrow{h \rightarrow 0} L[\phi]$ in L_{loc}^∞
- (Monotone) If $\psi(x) = \phi(x)$ and $\psi \geq \phi$ then $L_h[\psi](x) \geq L_h[\phi](x)$

but only for nondecreasing functions ϕ and ψ !

Numerical scheme

Given discretization parameters $h, \tau > 0$, a uniform grids $x_k = kh$ and $t_j = j\tau$:

$$V_k^{j+1} = V_k^j - \tau L_h V_k^j,$$
$$V_k^0 = \int_{-\infty}^{x_k} d\mu_0(x).$$

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Key result – Comparison for nondecreasing solutions: For V^j, W^j nondecreasing,

$$V^j \leq W^j \implies V^{j+1} \leq W^{j+1}.$$

provided CFL-type conditions:

$$\tau \leq h^{2s+m-1} \quad (V, W \text{ bounded}) \quad \text{or} \quad \tau \leq h^{\max\{1, 2s\}} \quad (V, W \text{ Lipschitz})$$

Theorem

Let $\{V_i^0\}_{i \in \mathbb{Z}}$ and $\{W_i^0\}_{i \in \mathbb{Z}}$ are nondecreasing and satisfying one of the following:

$$|V_i^0|, |W_i^0| \leq M, \quad i \in \mathbb{Z}, \quad \text{and} \quad \tau \leq C_1 h^{2s+m-1},$$

$$|V_i^0|, |W_i^0| \leq M, \quad \frac{|V_{i+1}^0 - V_i^0|}{h}, \frac{|W_{i+1}^0 - W_i^0|}{h} \leq L, \quad i \in \mathbb{Z}, \quad \text{and} \quad \tau \leq C_2 h^{\max\{2s, 1\}},$$

Then

- (a) (Nondecreasing) $V_i^j \leq V_{i+1}^j$.
- (b) (ℓ^∞ -stability) $|V_i^j| \leq M$.
- (c) (Nonegative) If $V_i^0 \geq 0$ then $V_i^j \geq 0$.
- (d) (ℓ^∞ -contraction) $\sup_{i \in \mathbb{Z}, j \in \mathbb{N}} |W_i^j - V_i^j| \leq \sup_{i \in \mathbb{Z}} |W_i^0 - V_i^0|$.
- (e) (Lipschitz-stability) $\sup_{i \in \mathbb{Z}, j \in \mathbb{N}} \frac{|V_{i+1}^j - V_i^j|}{h} \leq \sup_{i \in \mathbb{Z}} \frac{|V_{i+1}^0 - V_i^0|}{h}$.

Let \bar{V} be an interpolation of V (p.w. linear in space and p.w. constant in time).

Theorem

(a) (Continuous case) If $v_0 \in BUC(\mathbb{R}^d)$ then:

- (i) (Convergence) $\bar{V}_h \xrightarrow{h \rightarrow 0} v$ locally uniformly in \bar{Q}_T .
- (ii) (Limit) $v \in BUC(\mathbb{R}^d)$ is a viscosity solution of (IP).

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(b) (Jump discontinuity) If

$$v_0(x) = \begin{cases} 0 & \text{if } x < a, \\ M & \text{if } x \geq a. \end{cases}$$

(i) (Convergence) $\bar{V}_h \xrightarrow{h \rightarrow 0} v$ locally uniformly in $\bar{Q}_T \setminus \{(a, 0)\}$.

(ii) (Limit) $v \in C_b(\bar{Q}_T \setminus \{(a, 0)\})$ is a **discontinuous** viscosity solution of (IP).

Convergence by *half-relaxed limit method* (Barles, Perthame, Souganidis).

[1] Uniform-in- h L^∞ bounds. Then

$$\underline{v}(x, t) := \liminf_{(y, \rho, h) \rightarrow (x, t, 0)} \bar{V}_h(y, s) \quad \bar{v}(x, t) := \limsup_{(y, \rho, h) \rightarrow (x, t, 0)} \bar{V}_h(y, s)$$

and $\underline{v} \leq \bar{v}$.

[2] Monotonicity and consistency. Then

$$\begin{aligned} \underline{v} & \text{ discontinuous viscosity supersolution} \\ \bar{v} & \text{ discontinuous viscosity subsolution} \end{aligned}$$

[3] Comparison principle for (IP). Then

$$\bar{v} \leq \underline{v}$$

[4] All together:

$$\bar{v} = \underline{v} =: v = \lim_{h \rightarrow 0} \bar{V}_h$$

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Once a numerical method for (IP) is provided, the approximation for (FPE) is naturally given by numerical differentiation

$$U_i^j = \frac{V_{i+1}^j - V_i^k}{h}$$

and \bar{U} the corresponding p.w. constant interpolation.

Rubinstein-Kantorovich metric

$$d_0(f_1, f_2) = \sup_{\varphi \in \text{Lip}_{1,1}(\mathbb{R})} \int_{\mathbb{R}} (f_1(x) - f_2(x))\varphi(x)dx,$$

where $\text{Lip}_{1,1}(\mathbb{R}) = \{\varphi \in C(\mathbb{R}) : \|\varphi\|_{L^\infty} \leq 1, \|D\varphi\|_{L^\infty} \leq 1\}$ and D is the weak derivative.

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Theorem

$\mu_0 \in \mathcal{M}_+(\mathbb{R})$ and u the weak solution of (FPE).

- a** (L^1 case) Assume $\mu_0 \in L^1(\mathbb{R})$ (resp. $\mu_0 \in L^1(\mathbb{R}) \cap L^\infty(\mathbb{R})$), then we have

$$\sup_{t \in [0, T]} d_0(\bar{U}_h(\cdot, t), u(\cdot, t)) \xrightarrow{h \rightarrow 0} 0.$$

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where $\text{Lip}_{1,1}(\mathbb{R}) = \{\varphi \in C(\mathbb{R}) : \|\varphi\|_{L^\infty} \leq 1, \|D\varphi\|_{L^\infty} \leq 1\}$ and D is the weak derivative.

Theorem

$\mu_0 \in \mathcal{M}_+(\mathbb{R})$ and u the weak solution of (FPE).

- a** (L^1 case) Assume $\mu_0 \in L^1(\mathbb{R})$ (resp. $\mu_0 \in L^1(\mathbb{R}) \cap L^\infty(\mathbb{R})$), then we have

$$\sup_{t \in [0, T]} d_0(\bar{U}_h(\cdot, t), u(\cdot, t)) \xrightarrow{h \rightarrow 0} 0.$$

- b** (Fundamental solution) Assume $\mu_0 = M\delta_a$ then we have

$$\sup_{t \in [t_0, T]} d_0(\bar{U}_h(\cdot, t), u(\cdot, t)) \xrightarrow{h \rightarrow 0} 0.$$

Sketch of the proof:

- [Previous theorems] $\bar{V}(x, t) \rightarrow v(x, t) = \int_{-\infty}^x u(y, t) dy$ locally uniformly.
- [Interpolation choice] $\bar{V}(x, t) = \int_{-\infty}^x \bar{U}(y, t) dy$
- [Convergence I] For $\varphi \in \text{Lip}_{1,1}$ and $\partial_x \varphi$ with compact support:

$$\int \varphi(\bar{U} - u) dx = \int \varphi_x \int_{-\infty}^x (\bar{U} - u) dy dx = \int \varphi_x (\bar{V} - v) dx \xrightarrow{h \rightarrow 0} 0.$$

- [Equi-Tightness]

$$\int \bar{U}(1 - \chi_R) dx < \varepsilon, \quad \text{and} \quad \int \bar{u}(1 - \chi_R) dx < \varepsilon.$$

- [Convergence II] For $\varphi \in \text{Lip}_{1,1}$:

$$\int \varphi(\bar{U} - u) dx = \int \varphi \chi_R (\bar{U} - u) dx + \int \varphi (1 - \chi_R) (\bar{U} - u) dx \xrightarrow{h \rightarrow 0} 0.$$

Final comments

- The numerical method works in this specific framework just because:
 - $u \geq 0 \implies v$ nondecreasing.
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- We can cover measure data μ_0 (i.e. fundamental solution), since in the integrated variable it turns into a theory of discontinuous viscosity solutions.
- Numerical methods in dimensions $d > 1$?
- Numerical methods for sign-changing solutions?

Thank you

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