

# Pursuit-evasion dynamics in predator prey models.

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Workshop - Diffusion in Warsaw , 11/05/2023

# Outline:

- General mathematical framework: quasilinear system of reaction-diffusion equations
- Typical ecological model - prey-predator interactions
- Pursuit-evasion models
- Chemical signalling
- Direct/indirect taxis & pursuit-evasion models
- Linearisation at the space-homogeneous steady state and space-time patterns formation
- Numerical solutions
- Blow-up versus existence of global-in-time solutions

The talk is based on the joint papers with the ERCIM scholar Purnendu Mishra (at present in Norwegian University of Life Science)

- 1 Purnendu Mishra, D.W. *Repulsive chemotaxis and predator evasion in predator prey models with diffusion and prey taxis* Math. Models. Methods. Appl. Sciences (M3AS) (2022)
- 2 Purnendu Mishra, D.W. *Pursuit-evasion dynamics for Bazykin-type predator-prey model with indirect predator taxis*, J.D.E. (2023)

# Quasilinear parabolic system

- System of  $N$  quasilinear reaction-diffusion equations :

$$u_t = \nabla \cdot (A(u)\nabla u) + f(u) \quad \text{in } \Omega_T := \Omega \times (0, T),$$

where  $u : \Omega_T \mapsto \mathbb{R}^N$  and  $f : \mathbb{R}^N \mapsto \mathbb{R}^N$  with BC & IC.

- diffusion matrix (for  $r \neq s$  **cross-diffusion terms** )

$$A(u) = [a(u)^{r,s}]_{1 \leq r, s \leq N}$$

- Example: the case  $N = 2$

$$u_{1,t} = \nabla \cdot ((a(u)^{1,1}\nabla u_1 + (a(u)^{1,2}\nabla u_2) + f_1(u_1, u_2),$$

$$u_{2,t} = \nabla \cdot ((a(u)^{2,1}\nabla u_1 + (a(u)^{2,2}\nabla u_2) + f_2(u_1, u_2).$$

setting  $a^{1,1} = d_1 Id$ ,  $a^{1,2} = -u_1 \xi Id$ ,  $a^{2,1} = \chi u_2 Id$ ,  $a^{2,2} = d_2 Id$  we obtain

$$a(u) = \begin{pmatrix} d_1 & -\chi u_1 \\ \xi u_2 & d_2 \end{pmatrix} Id.$$

with  $d_1, d_2, \chi, \xi > 0$  and finally:

$$u_{1,t} = d_1 \Delta u_1 - \xi \nabla \cdot u_1 \nabla u_2 + f_1(u_1, u_2),$$

$$u_{2,t} = d_2 \Delta u_2 + \chi \nabla \cdot u_2 \nabla u_1 + f_2(u_1, u_2).$$

# Quasilinear parabolic system II

## Theorem

Suppose that :

- initial conditions  $u_{j,0} \in W^{1,p}(\Omega)$ ,  $p > n$  are non-negative functions,
- $\langle \nabla u_j, \nu \rangle = 0$  where  $\nu$  is the outer normal vector on smooth boundary  $\partial\Omega$ ,
- $a^{r,s}$ ,  $f_i = u_i \tilde{f}(u)$  are  $C^\infty$  smooth functions,
- $\nabla \cdot (A(u)\nabla u)$  is normally elliptic.

Then there exists  $T_{max} > 0$  such that there exists the unique local non-negative classical solution  $u$  defined in  $\Omega \times (0, T_{max})$ . It satisfies the boundary and initial conditions and

$$u \in (C([0, T_{max}] : W^{1,p}(\Omega)) \cap C^\infty(\bar{\Omega} \times (0, T_{max})))^N.$$

Moreover, if  $A(u) = [a(u)^{r,s}]_{1 \leq r,s \leq N}$  is a triangular matrix then either

$$\lim_{t \rightarrow T_{max}} \|u(t)\|_\infty = +\infty \quad \text{or} \quad T_{max} = \infty.$$

**Basic problems are located about the question:**

How does the interplay between  $f(u)$  and  $A(u)$  impact the properties of solutions ?

- Existence of global classical solutions versus blow-up of solution in finite time  
The prototype for the case of single semilinear equation is the Fujita problem (1966)

$$u_t = \Delta u + u^q \quad \text{in } \mathbb{R}^n \times (0, +\infty). \quad u(\cdot, 0) \geq 0.$$

- pattern formation - bifurcations from the constant steady state,
- existence of global-in-time weak solutions.

# The Keller-Segel model of chemotaxis

Patlak (1953) - Keller-Segel model (1972)

$W = W(x, t)$  density of some chemical released by the members of with density  $N(x, t)$ ,  
 $x \in \Omega \subset \mathbb{R}^n$  with smooth boundary

$\chi$ -chemotactic sensitivity parameter

$$\left\{ \begin{array}{l} N_t = D_N \Delta N + \gamma N - \nabla \cdot (\chi N \nabla W) \\ W_t = D_W \Delta W + \gamma N - \mu W \\ \text{with homogeneous Neumann boundary condition} \\ \langle \nabla N, \nu \rangle = \langle \nabla W, \nu \rangle = 0, \quad \text{on } \partial\Omega, \quad t > 0. \end{array} \right.$$

- (−) chemoattractant    (+) chemorepellent
- Early stages of the fruit body formation in slime mold *Dictostelium Discoideum* )
- For  $n = 1$  - global in time classical solution (Nagai, 1995)
- For  $n = 2$  - global solution for  $\int_{\Omega} N_0 dx$  small enough, otherwise  $T_{max} < \infty$ . (Nagai, Senba, Yosida; 1997, Biler, 1998)

# A typical prey-predator model (O.D.E. case)

$N(t)$ —prey density,

$P(t)$ —predator density

$$\frac{dP}{dt} = bF(N)P - \delta P := R_P(N, P),$$

$$\frac{dN}{dt} = rN \left( 1 - \frac{N}{K} \right) - F(N)P := R_N(N, P)$$

$F = F(N)$ -functional response e.g. -amount of food (prey) consumed per predator per unite of time, Holling's type II function:

$$F = F_H(N) = \frac{aN}{1 + T_h aN} \quad a, b > 0,$$

- 1 The Rosenzweig-MacArthur prey-predator model (1963)  
 $r$  - growth rate,  $\delta$  - death rate,  $a$ -attack rate,  $T_h$ - handling time.

For  $K = \infty$  and  $T_h = 0$  we get the Lotka-Volterra model.

$b$ — efficiency of conversion of food into offspring

- 2 For some set of parameters there is a unique globally stable steady state which may lose stability and limit cycle emerges via the Hopf bifurcation

# Chemical signalling

- Many chemicals (e.g. pheromones, kairomones) released by animals are used as means of inter and intraspecific communication - (chemical signaling) and sense of smell is a primary means by which prey animals detect predators or prey and trigger suitable behavioral responses .
- The chemical signal may be released by predator/prey itself (odor of predator or prey) or it may be released due to damage of prey captured (e.g. blood in aquatic ecosystems).
- Let  $W$  be a chemical released by prey or predator then the corresponding equation reads

$$W_t = d_3 \Delta W + g(N, P) - \mu W$$

where  $g = g(N, P)$  is the rate of chemical signal production and  $\mu$  is the degradation rate

$$g(N, P) = \gamma P \quad \text{or} \quad g(N, P) = \gamma N \quad \text{or} \quad g(N, P) = \beta F(N)P ,$$

# Terminology: Direct/indirect prey taxis and/or predator taxis

- direct prey-taxis is a directed movement of predator toward the gradient of prey density,
- **indirect prey-taxis** is a directed movement of predator toward the density gradient of a chemical released by prey,
- direct repulsive predator taxis is the directed movement of prey in the opposite direction to the gradient of predator density.
- **indirect repulsive predator taxis** is a directed movement of prey in the opposite direction to the density gradient of a chemical released by predator.
- **pursuit- evasion** model includes both direct/indirect prey taxis (**pursuit**) and repulsive direct/indirect predator taxis (**evasion**).
  
- In the context of predator-prey models the term indirect taxis was first used for a simplistic model in J.I. Tello, D.W. (M3AS, 2016).
- Similar idea was also used in in a different context in K. Fujie, T.Senba, (JDE, 2017)
- Tao, M. Winkler (J.Eur.Math. Soc., 2017)

# The prey-predator model with prey taxis (direct)

$$P_t = d_P \Delta P - \xi \nabla \cdot P \nabla N + bF(N)P - \delta P,$$

$$N_t = d_N \Delta N + rN \left(1 - \frac{N}{K}\right) - F(N)P.$$

with homogeneous Neumann boundary conditions (no-flux) and initial conditions

on smooth boundary  $\partial\Omega$ ,  $\Omega \subset \mathbb{R}^n$  and initial conditions. ( $\xi > 0$ )

- P.Kareiva, G.T. Odel (*Am. Naturalist* 1987),
- **Prey-taxis was found to stabilize prey-predator interactions**, no pattern formation is possible if ( $\xi > 0!$ )-J.M. Lee, T. Hillen and M.A. Lewis (*J. Biol. Dyn.*, 2009)

**Global-in-time existence of solutions:**

- $n \geq 1$  (with volume filling effect) B. Ainseba, et.al.(*NARWA*, 2008), Y. Tao (*NARWA*, 2010)
- $n \geq 1$  (classical sol., for small  $\xi$  with  $F(N)$  bounded) - S. Wu, J.Shi, B.Wu (*JDE* 2016); D.Li (*DCDS* 2021)
- $n \leq 2$  (**classical sol.**)- H.-Y Jin, Z.Wang (*JDE*, 2017), T. Xiang (*Nonlin Anal*, 2018), D. Li (*DCDS*, 2021)
- $n \leq 5$  (**weak solutions**) M. Winkler (*JDE*, 2017)

# Pursuit-evasion predator-prey model with direct taxis

$$\begin{cases} P_t = d_P \Delta P - \xi \nabla \cdot P \nabla N + R_P(P, N), \\ N_t = d_N \Delta N + \chi \nabla \cdot N \nabla P + R_N(P, N), \end{cases}$$

with homogeneous Neumann boundary conditions (no flux)

- **The main part of the system is not upper triangular** (full cross diffusion system)
- Formal stability/instability analysis, travelling waves)- Y. Tyutyunov, L. Titova, R.Arditi (Math. Mod.. Nat. Phenom., 2007)  
Global-in-time existence of solutions
- $n \leq 3$ - (class. sol. in a neighbourhood of the constant steady state)  
M. Fuest (SIAM J. Math. Anal, 2020)
- $n = 1$  - ( no restriction on the size of initial data, approximation by 6-th order operators)  
Y.Tao, M. Winkler (J.F.A, 2021) , (Nonlinear Anal. RWA, 2022)

# Pursuit evasion model - indirect taxis for both prey and predator

$$P_t = d_P \Delta P - \chi \nabla \cdot (P \nabla U) + R_P(P, N) - \delta_1 P^2,$$

$$N_t = d_N \Delta N + \xi \nabla \cdot (N \nabla W) + R_N(P, N) - r_1 N^2,$$

$$W_t = d_W \Delta W + \alpha_w P - \mu_w W,$$

$$U_t = d_U \Delta U + \alpha_u N - \mu_U U,$$

with homogeneous Neumann boundary conditions (no-flux)

- **The main part of the system is upper triangular**

**Global-in-time existence of solutions:**

- $n \leq 3$  (with  $\chi$  and  $\xi$  small enough or  $\delta_1, r_1$  big enough) - S. Wu (JMAA, 2022)

# Pursuit -evasion prey-predator model with indirect repulsive predator taxis and prey taxis

$$P_t = D_P \Delta P - \nabla \cdot (\xi P \nabla N) - \delta P + bF(N)P,$$

$$N_t = D_N \Delta N + \nabla \cdot (\chi N \nabla W) + rN \left(1 - \frac{N}{K}\right) - F(N)P,$$

$$W_t = D_W \Delta W + \gamma P - \mu W.$$

- Model B : ( $\chi > 0$   $\xi = 0$ ) **indirect repulsive predator taxis**
- Model A : ( $\chi > 0$   $\xi > 0$ ) **pursuit-evasion model**
- Basic  $L^1(\Omega)$  estimate :

$$\frac{d}{dt} \left( \int_{\Omega} P(x, t) dx + b \int_{\Omega} N(x, t) dx \right) + C_1 \left( \int_{\Omega} P(x, t) dx + b \int_{\Omega} N(x, t) dx \right) \leq C_2$$

where  $C_1$  and  $C_2$  are positive constants.

## Theorem

Suppose that  $P_0, N_0, W_0 \in W^{1,r}(\Omega)$ ,  $r > n$  are non-negative functions.  
For Model A and Model B there exists the unique local non-negative classical solution  $(N, P, W)$  satisfying boundary and initial defined on  $\bar{\Omega} \times [0, T_{max})$  such that

$$(N, P, W) \in (C([0, T_{max}) : W^{1,r}(\Omega)) \cap C^{2,1}(\bar{\Omega} \times (0, T_{max})))^3.$$

- Moreover,  $T_{max} = \infty$  and the solution is uniformly  $L^\infty$ -bounded in the case of
- Model B ( $\chi > 0, \xi = 0$ ) for all  $n \geq 1$
- Model A ( $\chi > 0, \xi > 0$ ) in the case of  $n = 1$ .
- P. Mishra, D.W. (Math. Mod. & Methods in Appl. Sc. (M3AS), 2022)

# Linear stability analysis for Model B and Hopf bifurcation

- The coexistence steady state to Model B is of form

$$\bar{E} = (\bar{N}, \bar{P}, \bar{W}) \quad \text{where} \quad \bar{W} = \frac{\mu}{\gamma} \bar{P}.$$

- A complex number belongs to the spectrum of the linearization of Model B at  $\bar{E}$  iff it is an eigenvalue of the following stability matrix :

$$M_j = \begin{pmatrix} -D_1 h_j + a_{11} & a_{12} & -\chi \bar{N} h_j \\ a_{21} & -D_2 h_j + a_{22} & 0 \\ 0 & a_{32} & -D_3 h_j + a_{33} \end{pmatrix}.$$

where  $\{h_j\}_{j=0}^{\infty}$  denotes the eigenvalues of the Laplace operator  $-\Delta$  with homogeneous Neumann boundary condition and  $[a_{i,j}]$  is the Jacobian matrix for O.D.E. case.

$$a_{11} < 0, \quad a_{12} < 0, \quad a_{21} > 0, \quad a_{22} \leq 0, \quad a_{32} > 0, \quad a_{33} < 0.$$

- For any  $\chi > 0$  considered as a **bifurcation parameter**:  $\det M_j < 0$  and  $\text{tr} M_j < 0$ .

# Linear stability analysis for Model B and Hopf bifurcation

The dispersal equation of stability matrix  $M_j$  is following

$$\lambda^3 + \rho_j^{(1)}\lambda^2 + \rho_j^{(2)}\lambda + \rho_j^{(3)}(\chi) = 0$$

where

$$\begin{aligned}\rho_j^{(1)} &= -\text{tr}M_j = -(a_{11} + a_{22} + a_{33}) + (D_1 + D_2 + D_3)h_j, \\ &:= \alpha_0 + \alpha_1 h_j,\end{aligned}$$

$$\begin{aligned}\rho_j^{(2)} &= a_{11}a_{22} - a_{12}a_{21} + a_{11}a_{33} + a_{22}a_{33} \\ &\quad + h_j(-a_{22}D_1 - a_{33}D_1 - a_{11}D_2 - a_{22}D_3 - a_{11}D_3 - a_{33}D_2) \\ &\quad + h_j^2(D_1D_2 + D_1D_3 + D_2D_3) \\ &:= \beta_0 + \beta_1 h_j + \beta_2 h_j^2,\end{aligned}$$

$$\begin{aligned}\rho_j^{(3)}(\chi) &= -\det M_j = -a_{11}a_{22}a_{33} + a_{12}a_{21}a_{33} \\ &\quad + h_j(a_{22}a_{33}D_1 + a_{11}a_{22}D_3 - a_{12}a_{21}D_3 + a_{11}a_{33}d_2) \\ &\quad + h_j^2(-a_{22}D_1D_3 - a_{33}D_1D_2 - a_{11}D_2D_3) + D_1D_2D_3h_j^3 + \chi a_{21}a_{32}\bar{N}h_j, \\ &= (\gamma_0 + \gamma_1 h_j + \gamma_2 h_j^2 + \gamma_3 h_j^3) + \chi(\gamma_4 h_j) := \rho_j^{(3,1)} + \chi\rho_j^{(3,2)} > 0\end{aligned}$$

where we have denoted  $\rho_j^{(3)}(\chi) = \rho_j^{(3,1)} + \chi\rho_j^{(3,2)}$ . It can be checked that all coefficients  $\alpha_j, \beta_j, \gamma_j$  are positive.

# Linear stability analysis for Model B and Hopf bifurcation

- 1  $\bar{E}$  is linearly stable if and only if for each  $j \geq 0$  matrices  $M_j$  have eigenvalues with negative real parts which according to the Routh-Hurwitz stability criterion is equivalent to the conditions

$$\rho_j^{(1)} > 0, \quad \rho_j^{(3)} > 0,$$

$$\text{and } Q_j := \rho_j^{(1)}\rho_j^{(2)} - \rho_j^{(3)}(\chi) = \rho_j^{(1)}\rho_j^{(2)} - \rho_j^{(3,1)} - \chi\rho_j^{(3,2)} > 0 \quad \text{for all } j \geq 0.$$

- 2 There exists  $\chi^H > 0$  such that

$$\chi^H = \min_{j \in \mathbb{N}_+} \tilde{\Psi}(h_j) := \left\{ \frac{\rho_j^{(1)}\rho_j^{(2)} - \rho_j^{(3,1)}}{\rho_j^{(3,2)}} \right\} \quad (1)$$

and the steady state  $\bar{E}$  is stable if  $\chi < \chi^H$ .

- 3 If

$$\tilde{\Psi}(h_j) \neq \tilde{\Psi}(h_k) \quad \text{for } j \neq k$$

then the minimum is attained for a single  $j = j_0$ .

- 4 Since  $\text{tr}M_{j_0} < 0$  and  $\det M_{j_0} < 0$  there is one real negative eigenvalue and a pair of conjugate eigenvalues which cross imaginary axis for  $\chi = \chi^H$  with the transversality condition being satisfied.

## Theorem

*There exist  $\chi^H > 0$  such that steady state  $\bar{E}$  in model B is locally asymptotically stable if  $\chi < \chi^H$ . Moreover, at  $\chi^H$  a solution periodic in space and time emerges according to the Hopf bifurcation mechanism.*

- 5 based on result of Amann, 1991

# Numerical solutions of models A and B

- Non-dimensional Rosenzweig-MacArthur model in the frame of model A

$$\begin{cases} N_t = \Delta N + \nabla \cdot (\chi N \nabla W) + rN(1 - N) - \frac{aN P}{(1 + \beta N)}, \\ P_t = d_p \Delta P - \nabla \cdot (\xi P \nabla N) - \delta P + \frac{cNP}{(1 + \beta N)}, \\ W_t = d_w \Delta W + \gamma P - \mu W, \end{cases}$$

with non-negative initial and no-flux boundary condition.

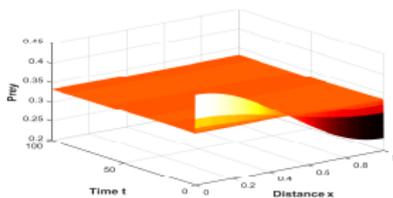
- 1D simulations with the help of MATLAB PDEPE tool ( $\Delta t = 0.01, \Delta x = 0.1$ )
- 2D simulations with the help of FreeFem++ solver ( $\Delta t = 0.01, \Delta x = \Delta y = 0.1$ )
- Values of model parameters are assumed to be

$$\begin{cases} r = 0.25, \beta = 2, c = 0.85, a = 0.95, \delta = 0.17, \\ \mu = 0.5, \gamma = 10, d_p = 0.01, d_w = 0.01. \end{cases} \quad (2)$$

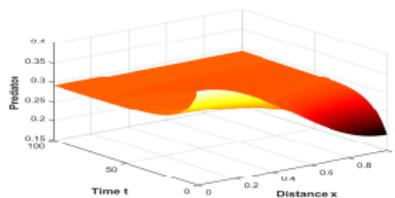
- Unique coexistence steady state  $E = (0.3333; 0.2924; 5.8490)$  and  $\chi^H = 6.889$ .
- Initial data : perturbation of the steady state e.g.

$$N(x, 0) = \bar{N} + 0.1 \cos\left(\frac{j\pi x}{L}\right), \quad P(x, 0) = \bar{P} + 0.1 \cos\left(\frac{j\pi x}{L}\right), \quad W(x, 0) = \bar{W} + \cos\left(\frac{j\pi x}{L}\right) \quad (3)$$

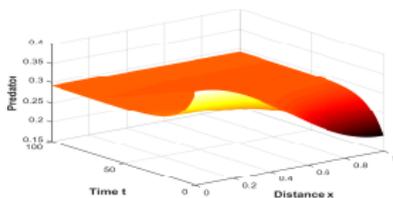
# Model B; convergence to the steady state (1D simulations)



(a)



(b)

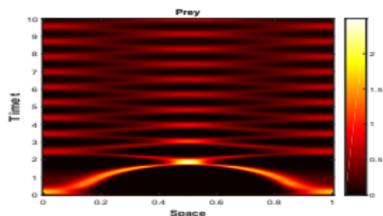


(c)

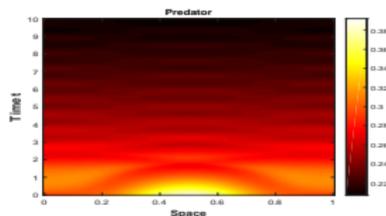
Figure 1: Model B: Perturbation in model B approaches the constant steady state  $\bar{E}$  for  $\chi < \chi^H$  with  $j = 1$

# Model B; transition of perturbation (1D simulation)

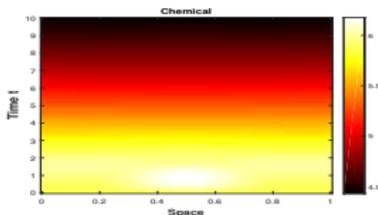
Initial data  $N(x, 0) = \bar{N}$ ,  $P(x, 0) = \bar{P} + 0.1e^{-\left(\frac{x-0.5}{0.2}\right)^2}$ ,  $W(x, 0) = \bar{W}$



(a)



(b)



(c)

Figure 2: Model B: spatio-temporal patterns for  $\chi > \chi^H$ .

# Model A; periodic solutions (1D simulations)

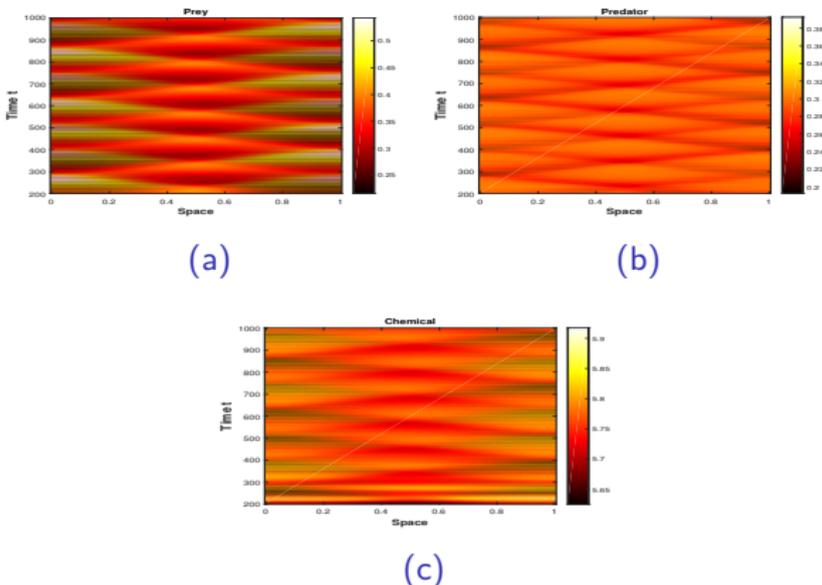
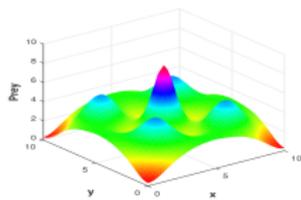


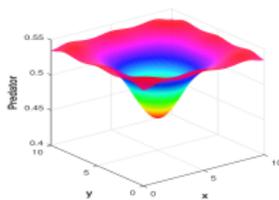
Figure 3: Model A: space-time patterns in unit domain when  $\chi = 5$ ,  $\xi = 0.2$  and symmetrical initial data with  $j = 4$ .

# Model B; separation regions (2D surface plot)

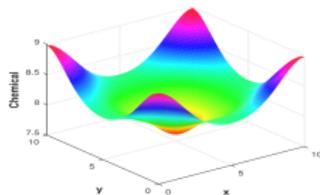
**Gaussian initial data** for predator centered in the middle of the square with constant initial data for the prey  $N = \bar{N}$  and for the chemical  $W = \bar{W}$



(a)



(b)

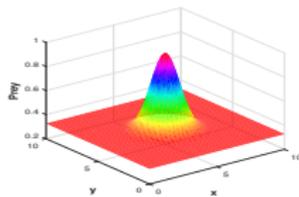


(c)

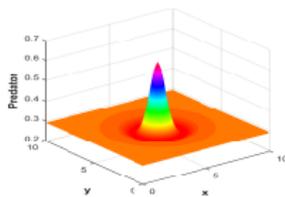
Figure 4: Model B ( $\xi = 0$ ): 2D separation regions for  $\chi = 10$  at time step  $t = 1500$ .

# Model A; spike solution

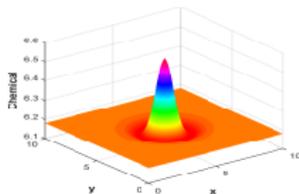
Gaussian initial data for predator and prey centered in the middle of the square with constant initial data  $\bar{W}$  for the chemical.



(a)



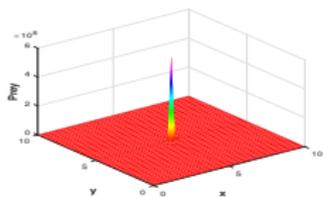
(b)



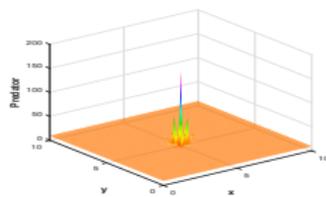
(c)

Figure 5: Model A: 2D simulation result for model A at time  $t = 10$  for  $\chi = 0.5, \xi = 10.0$

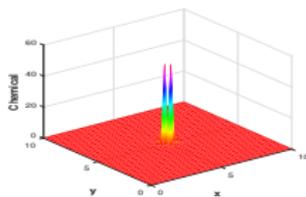
# Model A; spike solution



(a)



(b)



(c)

Figure 6: Model A: numerical indication of blowup at time  $t = 134$  for model A for  $\chi = 0.5$ ,  $\xi = 10.0$

# How to modify model A to prevent blow-up?—>

## Model C

- In Model C a minimal modification with respect to model A is made for prevention of blow-up in finite time.
- The kinetic part is as in the classical Bazykin model ( 1976).
- Density-dependent suppression of velocity in predators is interpreted as the result of interference (kind of regularisation)

$$\begin{cases} P_t = d_P \Delta P - \xi \nabla \cdot P \left( \frac{\nabla N}{1 + \sigma P} \right) + bF(N) - \delta P - \delta_1 P^2, \\ N_t = d_N \Delta N + \chi \nabla \cdot N \nabla W - F(N)P + rN - r_1 N^2, \\ W_t = d_W \Delta W + \gamma P - \mu W, \end{cases}$$

## Theorem

If  $P_0, N_0, W_0 \in W^{1,r}(\Omega)$ ,  $r > n$  are non-negative functions then there exist global in-time, non-negative classical solution to Model C satisfying boundary and initial condition provided  $n \leq 3$  and the following restrictions on parameters are satisfied

$$\mathbf{Q} \begin{cases} \delta_1 \geq \left( \frac{\gamma^2(16+n)}{d_W} + d_W \right), \\ r_1 \geq \left( \frac{\chi^2 A_N}{(d_N)^2} + \frac{2\chi^2}{d_W} + d_W \right), \\ \text{with } A_N = \frac{2 \left( (d_N)^2 + (d_W)^2 + \xi^2 \sigma^{-2} \right)}{d_W}. \end{cases}$$

P.Mishra, D.W., JDE. 361 (2023)391-416 .

# Numerical solutions to Model C

- Set of parameters

$$r = 2, r_1 = 1.8, a = 0.7, b = 0.9, \beta = 2, \mu = 0.01, \delta = 0.1, \delta_1 = 0.15, \\ \gamma = 0.015, d_n = 1, d_p = 0.1, d_w = 0.05.$$

- For this choice of parameters values the restriction **Q** holds if and only if  $\sigma > \sigma_c := 19.7$
- For  $\sigma < \sigma_c$ , num. sol. to Model C exhibits finite-time blow-up of solution
- For  $\sigma > \sigma_c$  there is prevention of blow-up (global solutions) .
- Initial data: perturbation of the constant steady state  
 $E^* = (P^*, N^*, W^*) = (0.741, 1.016, 0.74)$

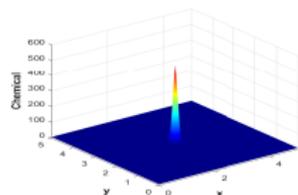
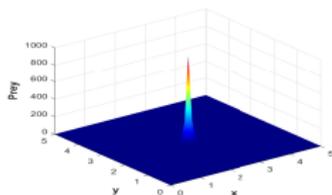
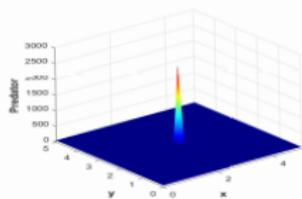
$$P_0(x, y) = P^* + 500e^{-100((x-2.5)^2+(y-2.5)^2)},$$

$$N_0(x, y) = N^* + 800e^{-100((x-2.5)^2+(y-2.5)^2)},$$

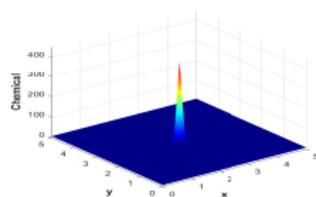
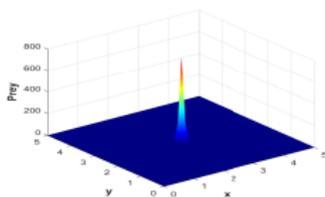
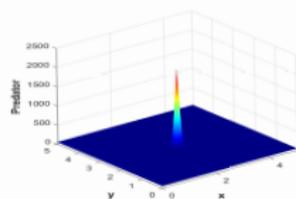
$$W_0(x, y) = W^* + 100e^{-100((x-2.5)^2+(y-2.5)^2)}$$

where  $(x, y) \in \Omega = (0, 5) \times (0, 5)$

# Figure 1

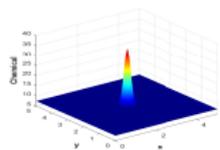
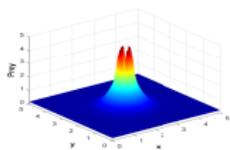
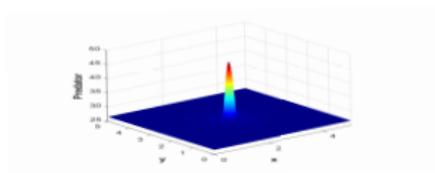


(a)

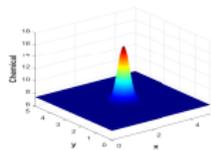
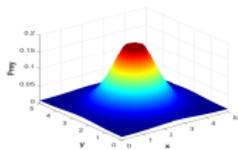
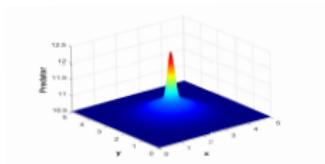


(b)

Figure 7: (a) Approximated blowup solution at time  $t = 1.5 \times 10^{-4}$  for  $\sigma = 0.0$  (b) Approximated blowup solution at time  $t = 2.3 \times 10^{-4}$  for  $\sigma = 5.0$  subject to initial conditions. It was assumed  $\chi = 0.1$ ,  $\xi = 30$ .

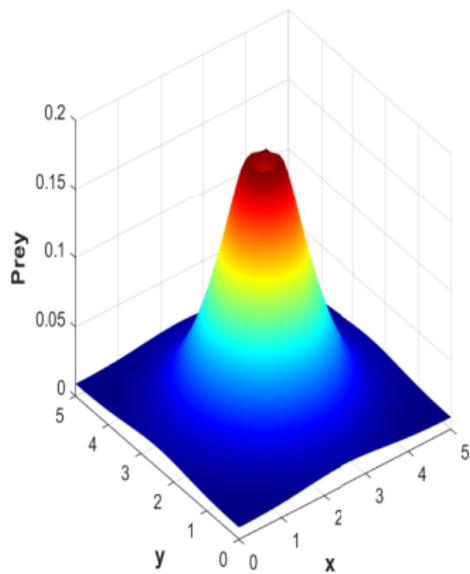


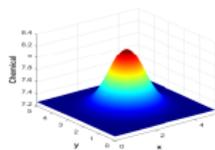
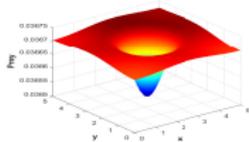
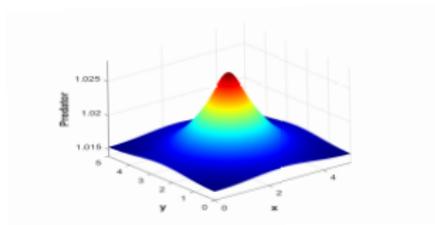
(a)



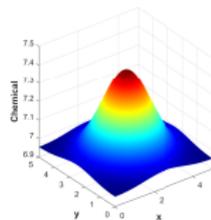
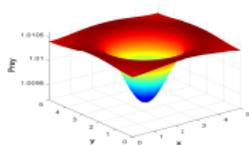
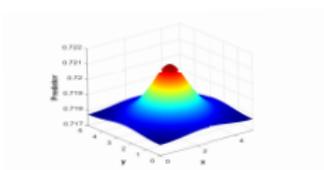
(b)

Figure 8: Snapshots for  $\sigma = 25$  at different time steps. (a)  $t = 13$ , (b)  $t = 50$





(a)



(b)

Figure 9: Snapshots for  $\sigma = 25$  at different time steps. (a)  $t = 100$ , (b)  $t = 500$ . All other parameter values and initial condition is same as in figure (7).

# Sketch of proof of the global existence for Model C

- 1 There is a local smooth solution defined on  $[\tau, T_{max})$  satisfying  $L^1(\Omega)$ -bound.
- 2 We begin with the N-equation

$$N_t = d_N \Delta N + \chi \nabla \cdot N \nabla W - F(N)P + rN - r_1 N^2$$

- 3 Using the Gagliardo-Nirenberg inequality and  $L^1(\Omega)$ -bound one proves that for  $n \leq 3$

$$\sup_{t \in [\tau, T_{max})} \|N(\cdot, t)\|_k \leq C_N(k) \quad \text{for any } k \geq 1$$

provided

$$\sup_{t \in [\tau, T_{max})} \|\nabla W(\cdot, t)\|_4 \leq C'_W.$$

Then

$$\sup_{t \in [\tau, T_{max})} \|N(\cdot, t) \nabla W(\cdot, t)\|_{4-\varepsilon} \leq C''_W$$

- 4 Using properties of the heat semigroup we infer that

$$\sup_{t \in [\tau, T_{max})} \|N(\cdot, t)\|_\infty \leq C_N.$$

- 5 Using  $L^p - L^q$  estimates for analytic semigroups ( $n \leq 3$ ) we get

$$\sup_{t \in [\tau, T_{max})} \|\nabla N(\cdot, t)\|_p \leq C'_N \quad \text{for } p < 4$$

Next it is easy to deduce by parabolic regularity that

$$\sup_{t \in [\tau, T_{max})} \|\nabla P(\cdot, t)\|_\infty \leq C_P, \quad \sup_{t \in [\tau, T_{max})} \|\nabla W(\cdot, t)\|_\infty \leq C_W$$

# Sketch of proof of the global existence for Model C

- The most complicated part of the proof is to find estimate on  $\|\nabla W(\cdot, t)\|_4$ . To this end we derive differential inequality

$$y'(t) + y(t) \leq \text{Const.} \quad \text{for } t \in [\tau, T_{\max})$$

where for suitable constants  $A_1$  and  $A_2$

$$y(t) = \int_{\Omega} |\nabla W|^4 + \int_{\Omega} P |\nabla W|^2 + \int_{\Omega} N |\nabla W|^2 + A_1 \int_{\Omega} N^2 + A_2 \int_{\Omega} P^2.$$

- We use Bochner's type inequality : For  $W \in C^2(\bar{\Omega})$  there holds

$$2\nabla W \nabla \Delta W = \Delta |\nabla W|^2 - 2|D^2 W|^2$$

and

- Mizoguchi-Souplet inequality : for  $u \in C^2(\bar{\Omega})$  satisfying  $\frac{\partial u}{\partial \nu} = 0$  on  $\partial\Omega$  and  $\Omega$  there holds the following pointwise inequality

$$\frac{\partial |\nabla u|^2}{\partial \nu} \leq K |\nabla u|^2 \quad \text{on } \partial\Omega$$

where  $K$  depends on the curvature of  $\partial\Omega$ .

### Lemma

Let  $(P, N, W)$  be a solution to Model C. Then there exists a constant  $C > 0$  such that for  $t \in (0, T_{max})$ .

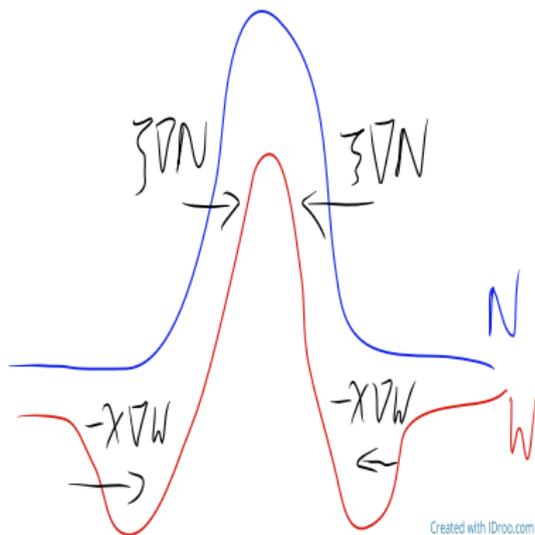
$$\begin{aligned} \frac{d}{dt} \int_{\Omega} |\nabla W|^4 + d_W \int_{\Omega} |\nabla(|\nabla W|^2)|^2 + 4\mu \int_{\Omega} |\nabla W|^4 \\ \leq \gamma^2 \left( \frac{16+n}{d_W} \right) \int_{\Omega} |\nabla W|^2 P^2 + C. \end{aligned}$$

# Corollaries and open questions

- Chemical signalling may destabilize a space-homogeneous steady state in a prey -predator model and gives rise to space-time dependent pattern formation.
- When an O.D.E. model is extended to a P.D.E model with taxis terms some mechanism of blow-up prevention might be necessary to be built in the model.
- None of the two taxis mechanisms studied in Model C alone can lead to the blow-up for  $n = 2$ . Their cumulative effect leading to blow-up demands farther investigation.
- Are there any weak solutions for for Model C when  $\sigma = 0$ , weak enough to grasp the singular solutions?

**Thank you.**

# The explanation of spiky solution formation



Cumulative effect of prey taxis and indirect predator taxis leads to aggregation .